

# The property of solutions for the equations of large-scale atmosphere with the non-stationary external forcings\*

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The long-term behavior of the atmospheric evolution, which cannot be answered and solved by the numerical experiments, must be understood before we design the numerical forecast models of the long-range weather and climate. It is necessary to carry out some studies of basic theory. Based on the stationary external forcings, Chou<sup>[1-3]</sup> studied the adjustment of the nonlinear atmospheric system tending to forcings in  $R^n$ . Then these results were extended to the infinite dimensional Hilbert space<sup>[4]</sup>. For the real atmospheric system, the external forcings are non-stationary. In such a case, are the results true or not? This note will deal with the main problem.

## 1 Problem, assumption and notations

The partial differential equations of large-scale atmosphere can be turned into the following operator equation:

$$\begin{cases} B \frac{\partial \psi}{\partial t} + (N+L)\psi = \xi, & (1) \\ B\psi = B\psi_0, \text{ when } t = t_0, & (2) \end{cases}$$

where  $B$  and  $L$  are the positively definite self-adjoint operators,  $N$  the anti-adjoint operator, both the non-homogeneous terms and the boundary conditions of the partial differential equations are included in  $\xi$  (for terms of  $B$ ,  $N$  and  $L$  see references [1-3, 5]).

In  $R^n$ ,  $B$  and  $L$  are the  $(n \times n)$  positively definite symmetric matrix,  $N$  the  $(n \times n)$  antisymmetric matrix. If we study the stationary or non-stationary external forcings,  $\xi$  is the  $n$ -dimensional constant vector or varied vector  $\xi_t$ , respectively. Our problem is the asymptotic property of atmospheric motion with the non-stationary external forcings as time  $t \rightarrow \infty$  in  $R^n$ , namely, the destination of the integral curve determined by the following equation as time  $t \rightarrow \infty$

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$$B \frac{\partial \psi}{\partial t} + (N+L)\psi = \xi_t. \quad (3)$$

We need an assumption and some notations before discussion. In reality, the external forcings should be bounded, namely

$$0 < \|\xi_t\|^2 \leq M < \infty, \quad (4)$$

where  $\|\cdot\|$  represents the norm.

For the sake of convenience, let  $M_0 = M(t)|_{t=t_0}$  represent one point of phase space. Then the solution to eq. (3) in which the initial value equals  $M_0$  is  $M(t) = \psi(M_0, t)$ , the following point set

$$\gamma = \{\psi(M_0, t) : t_0 \leq t < \infty\} \text{ for a certain initial value } M_0 \quad (5)$$

represents paths through  $M_0$ , and the following set of points

$$V(t^*) = \{\psi(M_0, t) : M_0 \in V_0\}_{t=t^*} \quad (6)$$

may be regarded as a mapping. It maps  $V_0$  to  $V$ , where  $V_0$  and  $V \subset \mathbb{R}^n$ .

## 2 Main theorems

**Lemma 1.** If the ellipsoid  $E: \sum_{i=1}^n (\psi_i - a_i)^2 / r_i = 1$ , where  $r_i$  and  $a_i$  ( $i=1, 2, \dots, n$ ) are constants,  $r_i \in \mathbb{R}^+$ ,  $a_i \in \mathbb{R}$ , then

i) for any point  $\psi$  out of  $E$ , we have  $\sum_{i=1}^n (\psi_i - a_i)^2 / r_i > 1$ ;

ii) for any point  $\psi$  inside of  $E$ , we have  $\sum_{i=1}^n (\psi_i - a_i)^2 / r_i < 1$ .

**Lemma 2.** If the ellipsoid  $E(A, r): \sum_{i=1}^n (\psi_i - a_i)^2 / r_i = 1$ , and  $A \in \mathbb{R}_1^n$ ,  $r \in \mathbb{R}_2^n$ ,  $r_i \in \mathbb{R}^+$  ( $i=1, 2, \dots, n$ ), where  $A = (a_1, a_2, \dots, a_n)$ ,  $r = (r_1, r_2, \dots, r_n)$ ,  $\mathbb{R}_1^n$  and  $\mathbb{R}_2^n$  are the bounded closed set, then

i)  $E = \bigcup_{R_1^n} E_{(A, r)}$  is a bounded set, where  $\bigcup_{R_1^n} E_{(A, r)}$  represents the union of all ellipsoids that take  $A$  as the center of ellipsoid and  $r$  as the axis,  $A$  is any point in  $\mathbb{R}_1^n$ ;

ii) for any point  $\psi$  out of  $E$ ,  $\sum_{i=1}^n (\psi_i - a_i)^2 / r_i > 1$ .

**Theorem.** Under the assumption of eq. (4), the solution to eq. (3) is satisfied, there exists a bounded closed set  $V_0$  such that

i) if  $M_0 \in V_0$ , then  $\gamma \in V_0$ ;

ii)  $\forall M_0 \notin V_0$ , there exists a  $\tau > 0$ , the point set  $V \in V_0$ , where  $V_\tau = \{\psi(M_0, t) : \tau \leq t < \infty\}$ .

### 3 Proof of theorems

Dotting both sides of eq. (3) with  $\psi$  and using the relation  $(\psi, N\psi) = 0$ , we have

$$\frac{d(\psi, B\psi)}{dt} = 2[(\psi, \xi_D) - (\psi, L\psi)]. \quad (7)$$

Because  $L$  is the  $(n \times n)$  positively definite symmetric matrix, there are  $n$  characteristic values  $\lambda_i > 0$  ( $i=1, 2, \dots, n$ ). Let  $e_i$  ( $i=1, 2, \dots, n$ ) be the corresponding normalized characteristic vectors. They form a set of the normalized orthogonal basis of the  $n$ -dimensional space. So let

$$\xi_i = \xi_{i1}(t)e_1 + \xi_{i2}(t)e_2 + \dots + \xi_{in}(t)e_n, \quad (8)$$

$$a_i = a_{i1}(t)e_1 + a_{i2}(t)e_2 + \dots + a_{in}(t)e_n, \quad (9)$$

where

$$a_i(t) = \xi_i(t)/(2\lambda_i), \quad (10)$$

then

$$2La_i = \xi_i. \quad (11)$$

Let

$$\psi = \psi_1 e_1 + \psi_2 e_2 + \dots + \psi_n e_n. \quad (12)$$

Because

$$Le_i = \lambda_i e_i, \quad (13)$$

$$(\psi, \xi_D) - (\psi, L\psi) = \sum_{i=1}^n \lambda_i \{a_i^2(t) - [\psi_i(t) - \psi_i]^2\}. \quad (14)$$

Consider the point set  $E_1$ :  $(\psi, \xi_D) - (\psi, L\psi) = 0$ , namely, the point set satisfies the following equation:

$$\sum_{i=1}^n \left\{ [\psi_i - q_i(t)]^2 \left/ \left[ \sum_{j=1}^n \lambda_j a_j^2(t) / \lambda_j \right] \right. \right\} = 1. \quad (15)$$

According to eqs. (4) and (10), we have  $0 < \|a_i\|^2 \leq M_1 < \infty$ ,  $a_i \in R_1^n$ , where  $R_1^n$  is a bounded closed set. Meanwhile,  $\lambda_i$  ( $i=1, 2, \dots, n$ ) belong to  $R^+$  and are bounded, so  $0 < \|r_i\|^2 \leq M_2 < \infty$ ,

where  $r_i = (r_{i1}(t), r_{i2}(t), \dots, r_{in}(t))$ ,  $r_i(t) = \sum_{j=1}^n \lambda_j a_j^2(t) / \lambda_j$ ,  $r_i(t) \in R^+$ . Then we have  $r_i \subset R_2^n$ , where  $R_2^n$  is

a bounded closed set. Let  $E = \bigcup_{R_1^n} E_{(a_i, r_i)}$  represent the union of all ellipsoids that satisfy eq. (15) and take  $a_i$  as the center of ellipsoid and  $r_i$  as the axis, where  $a_i$  is any point in  $R_1^n$ . As a result of Lemma 2,  $E$  is a bounded closed set; besides,  $E \supset E_1$ . Then according to Lemma 2, for any point out of  $E$ , we have

$$\frac{d(\psi, B\psi)}{dt} < 0. \quad (16)$$

Let a bounded closed set  $V_0 \supset E$ . Then  $\psi$  satisfies eq. (16), where  $\psi$  is any point out of  $V_0$ . Then there exists a certain time  $\tau$ , paths through any point out of  $V_0$  will run into  $V_0$  as time  $t > \tau$ , and paths through any point inside of  $V_0$  cannot run away from  $V_0$ . It is evident that there exists a bounded closed set  $V_0$  that satisfies  $V_0 \supset E$ . Then the theorem has been proved.

#### 4 Discussion

According to the above discussion, for the non-stationary external forcings with condition (4), the atmospheric system will run into the attractive set of points  $V_0$  while the time  $t$  is greater than a certain critical time  $\tau$ . Points out of  $V_0$  have nothing to do with the asymptotic behavior for the time  $t \rightarrow \infty$  and have only transient sense. The long-term behavior of the system will depend on the bounded closed set  $V_0$ . The result shows that the long-range weather or climate is in a state of attractor. Because the contracted evolution of dissipative system from the high-dimensional phase space to the low-dimensional attractor is a process that merges degrees of freedom, the effective degrees of freedom that determine the long-term behavior of the system are eventually reduced to limited degrees. We can make an estimate of dimension of the large-scale atmospheric attractor with the non-stationary external forcings in theory using the property of long-term weather. It is a great help to build a new theory and computational methods of long-range numerical forecast. This will be reported in other papers.

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